The role of dense, molecular gas during early stages of galaxy formation and evolution

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with



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Kostas Tassis (U.Chicago) Much of the gas in dwarf ξ LSB galaxies (and outskirts of normal spirals) is inert to star formation NGC 2915 (Blue Compact Dwarf)





CO map: Blitz et al. 03; HI map: Deul & van der Hulst (1987)

Kennícutt-Schmídt relatíon ín M33 ís very steep

Heyer et al. 2004; see also Boissier et al. 2003



Dwarfs look like the outskirts of massive disks in terms of SF



Adopted from a talk by F. Walter

Making sense of it

Locally, in individual molecular clouds:

$$\dot{\rho}_{\star} = \frac{\rho_g}{t_{\star}}$$

Krumholz & Tan (2006) show that for dense, molecular gas t_* scales as the free fall time of the dense gas: $t_{\star} = t_{ff}/SFR_{ff}$ with SFRff approximately independent of gas density



When star formation in a disk is averaged on some (~kpc) scale:

$$\langle \dot{\rho_{\star}} \rangle \propto \langle f_{\rm sf} \rangle \langle \rho_g \rangle^{1.5}$$

fraction of total gas mass eligiible for star formation

$$\begin{split} \Sigma_{\rm SFR} \propto \langle \dot{\rho_{\star}} \rangle \, h_{\rm SFR} \\ \Sigma_g \propto \langle \rho_g \rangle \, h_g^{\text{scale-height}} \, {}_{\rm of \ young \ stars} \end{split}$$

scale-height of all gas

$$\Sigma_{\rm SFR} \propto \langle f_{\rm SF} \rangle \frac{h_{\rm SFR}}{h_g^{1.5}} \Sigma_g^{1.5}$$

An example: H2-based star formation in a model disk galaxy $\dot{\rho}_{\star} = f_{\mathrm{H2}} \frac{\rho_g}{t}$ $-t_{\star} = 0.7 \mathrm{Gyr}$ $f_{\rm H2} = f_{\rm H2} \left(\rho_g, T, Z_g, U_{\rm isrf} \right)$ Robertson & Kravtsov, ApJ submitted (astro-ph/0710.2102) H D-SF+ISBF T = 0.3T = 0.3H D-SF+ISRF T = 0.3H D-SF+ISRF 2 =300 km/s 2 =125 km/s $v^2 = 50 \text{ km/s}$ gas surface density (colored by T) massive galaxy intermediate galaxy dwarf galaxy 10¹ SFR surface density H_-Cooling+ISRF H_-Cooling+ISRF H -Cooling+ISRF $v^{2} = 300 \text{ km/s}$ $v^{2} = 125 \text{ km/s}$ $v^{2} = 50 \text{ km/s}$ = 2044= 3.277= 4.238yr⁻¹ kpc⁻²] 10 10^{-2} Σ_{str} [h² M _{sun} 10⁴ slope=3.3 slope=2 slope=4.2 10

 $\Sigma_{gas} [h M_{sun} pc^{-2}]$ total gas surface density

 10^{2}

10¹

 10^{3}

10¹

 $\Sigma_{\text{gas}} [h M_{\text{sun}} pc^{-2}]$

 10^{2}

 10^{3}

 10^{3}

10-5

 10^{1}

 10^{2}

 $\Sigma_{\text{gas}} [h M_{\text{sun}} pc^{-2}]$

Scaling with H2 fraction and scale-heights is as expected



Modelíng H2 ín galaxy formatíon símulatíons

dense, molecular gas traces densest, high-pressure regions of the ISM



Galaxy formation simulation (ART code) with approximate 3D radiative transfer and a model for H2 formation on dust with approximate self-shielding using Sobolev approximation face-on and edge-on views of HI and H2 distribution in a z~4 gas disk



Gnedin, N., Tassis, Kravtsov 2008, in prep

Molecular fraction as a function of gas surface and local 3D density



Strong trends with metallicity (dust content) and local UV flux

Gnedin, N., Tassis, Kravtsov 2008, in preparation

Implications

At high z (lower metallicities, higher UV flux), dense star forming gas is expected to be more compact compared to local galaxies. Most of the gas in low-mass systems, may then stay in atomic gas and be inert to star formation. This can have implications for a number of observations and theoretical expectations about galaxy evolution.

Steeper Kennicutt-Schmidt relations can be expected in high-z galaxies. This can explain why dense DLA systems do not show the expected associated UV flux, predicted by the local KS relation (Wolfe & Chen 2007, Wild et al. 2007). Also, *lower H2 content of DLAs* compared to the MW, given their metallicity (e.g.Noterdaeme et al. 2008)

 Existence of very dense, compact (re<1 kpc) galaxies at z>2-3 (e.g., Zirm et al. 2007, Toft et al. 2007)

Low escape fractions of ionizing UV photons from most high-z galaxies (Gnedin, Kravtsov & Chen 2008).

□ *Existence of undisturbed, massive stellar disks at z~2.0-3.0* (Stockton et al. 2007, Tacconi et al. 2008).

Distribution of young stars and HI in two simulated high-z galaxies



If young stars are deeply embedded in extended HI gas, the resulting escape fraction of UV photons is very small

